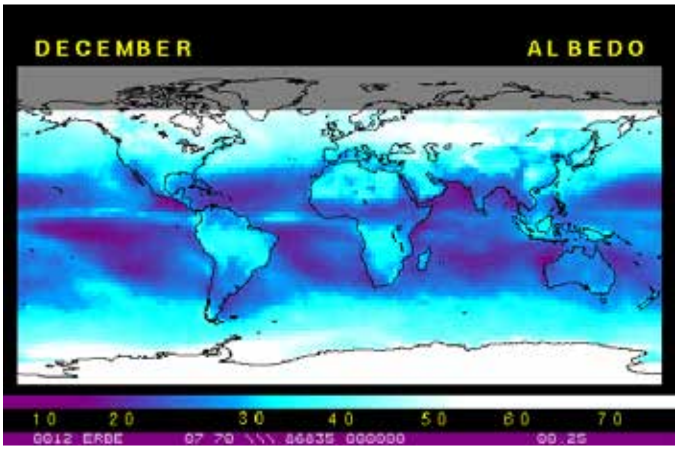
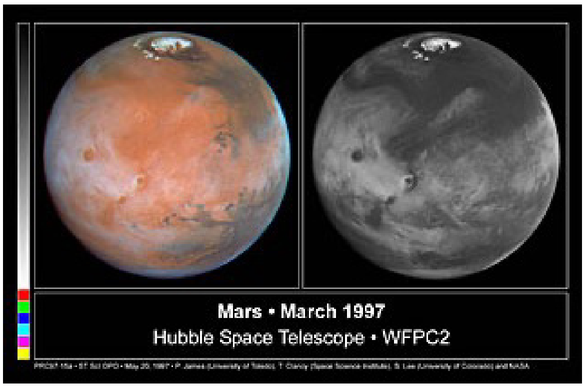
Atmospheric dynamics I

Clouds on the terrestrial planets



Earth
albedo 0.30

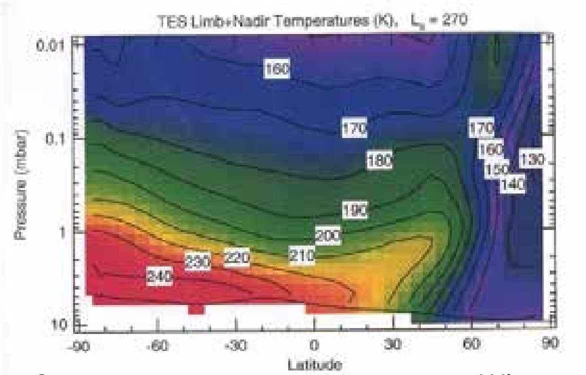


Mars
albedo 0.16

Venus
albedo 0.78

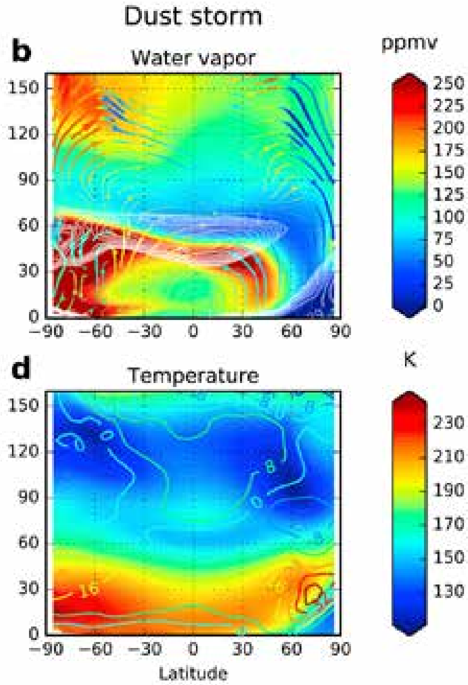
In one-dimensional model, the atmosphere is assumed to be horizontally uniform; however, this is not a good assumption in general.

Meridional distribution of zonal-mean temperature obtained by MGS/TES (Smith et al. 2001)

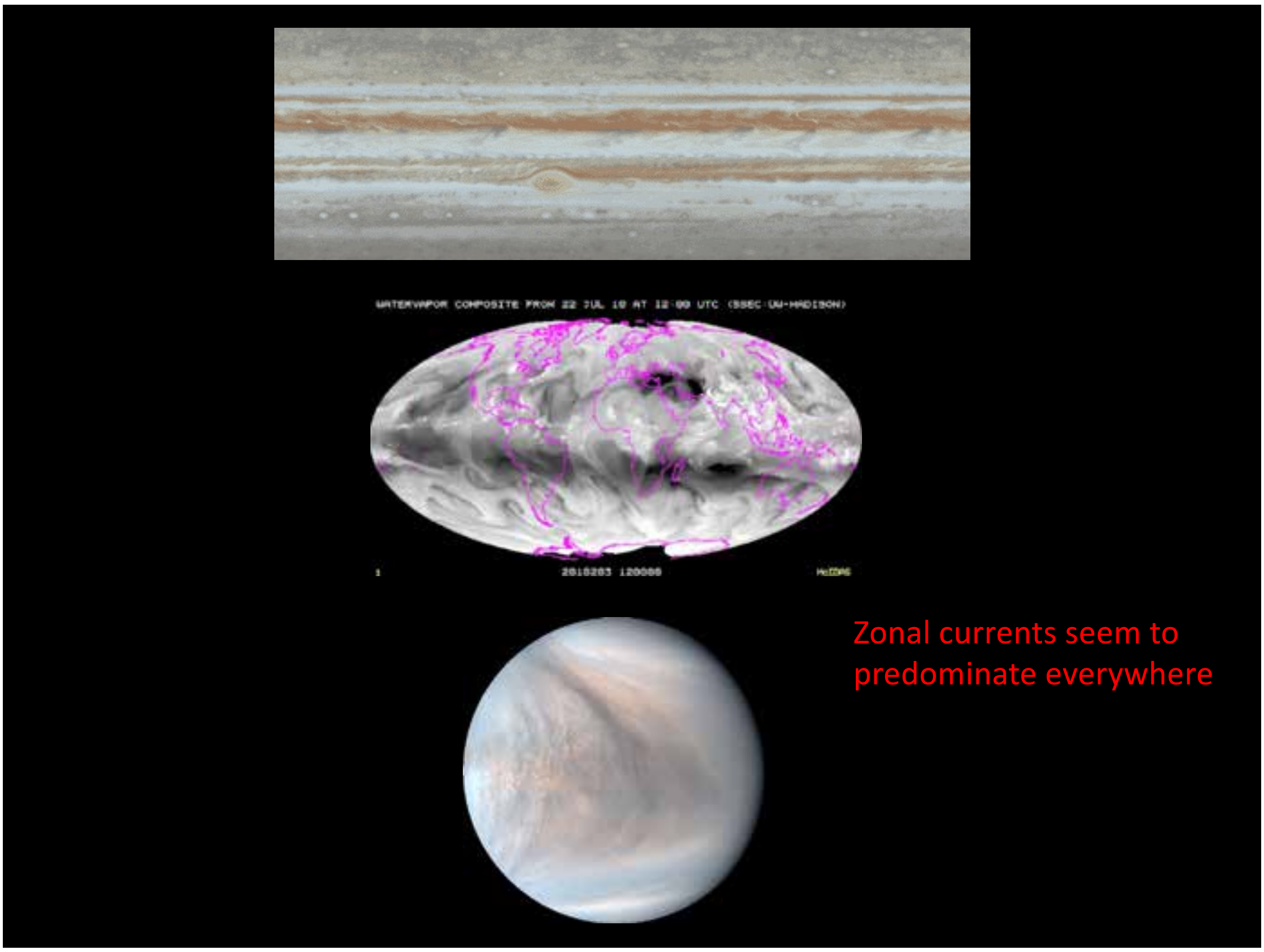


Summer Winter

Meridional cross sections in Mars GCM (Shaposhnikov et al. 2019)



3-D atmospheric circulation can play crucial roles in material transport



Parameter study of the atmospheric circulation of Earth-like planets with general circulation models (GCMs)

Williams (1988)

Meridional stream function

white : anti-clockwise

shade : clockwise

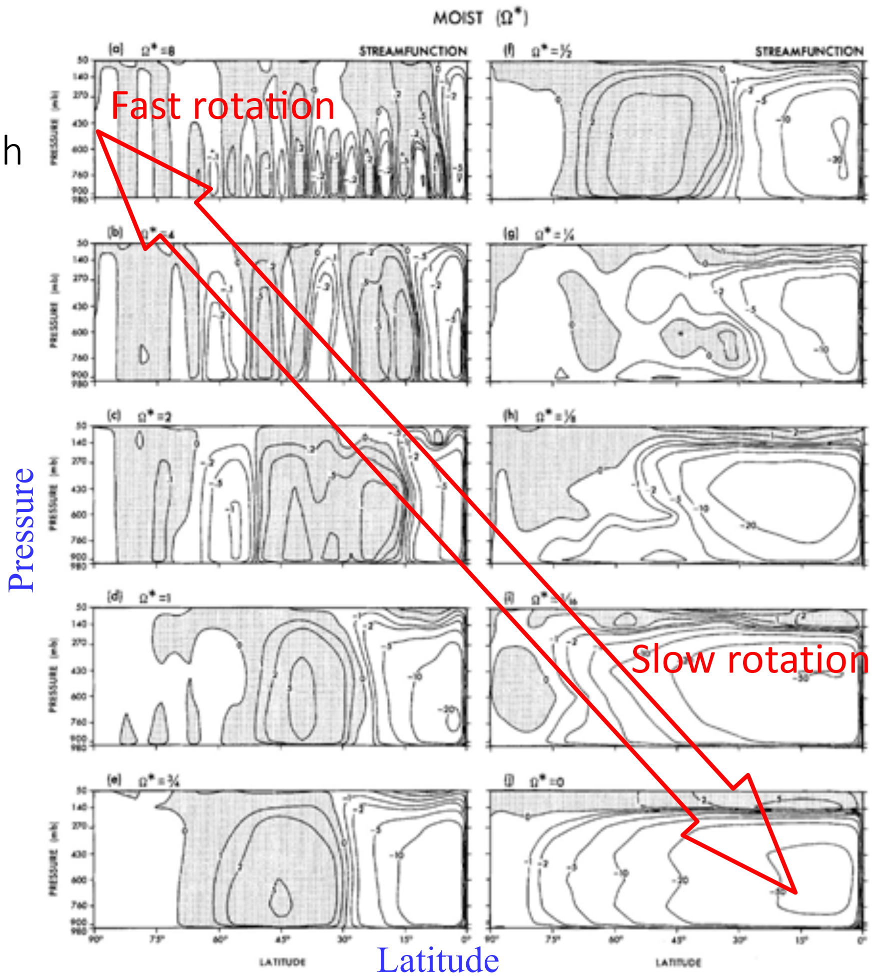
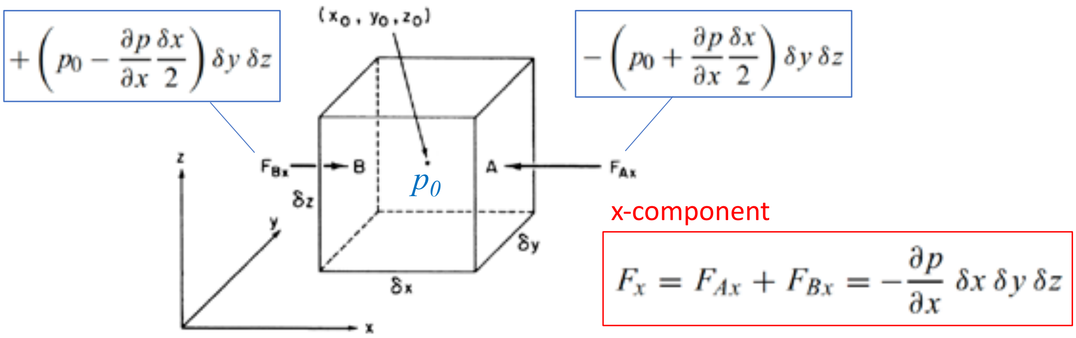


Fig. 3. Meridional distribution of the mean stream function for the MOIST model with Ω* = 0-8. Units: 10^{13} g s^{-1}

Momentum equation

Pressure gradient force:



The x component of the pressure gradient force acting on a fluid element.

Holton, 2004

Momentum equation in the inertial coordinate system

$$\frac{d\vec{v}}{dt} = -\frac{1}{\rho}\nabla p + \vec{g}$$

$\vec{v}$: velocity

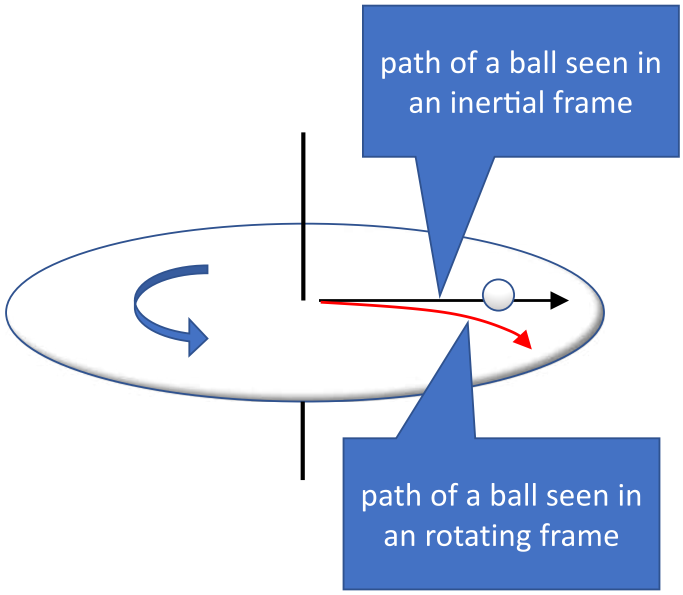
t : time

ρ : atmospheric density

p : atmospheric pressure

$\vec{g}$: gravity acceleration

Momentum equation in a rotating frame



Momentum equation in a rotating frame

Let us consider a frame rotating at an angular velocity $\vec{\Omega}$.

For a position vector $\vec{r}$, the relationship between the time derivative in the inertial frame d_r/dt and that in a rotating frame d/dt is

$$\frac{d\vec{r}}{dt} = \frac{d_r\vec{r}}{dt} + \vec{\Omega}\times\vec{r}$$

$$\therefore \vec{v} = \vec{v}_r + \vec{\Omega}\times\vec{r}$$

$\vec{v}_r$: velocity in rotating frame

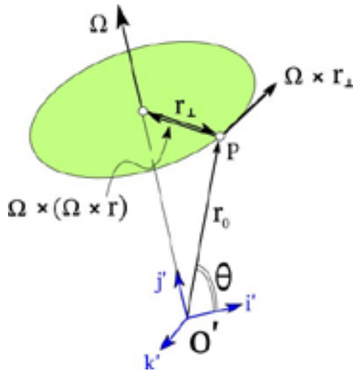
Similarly

$$\frac{d\vec{v}}{dt} = \frac{d_r\vec{v}}{dt} + \vec{\Omega}\times\vec{v}$$

Combining these, using $\vec{\Omega}\times(\vec{\Omega}\times\vec{r}) = -\Omega^2\vec{R}$ we have

$$\frac{d\vec{v}}{dt} = \frac{d_r\vec{v}_r}{dt} + 2\vec{\Omega}\times\vec{v}_r - \Omega^2\vec{R}$$

$\vec{R}$: vector from the rotational axis to the fluid position



Substituting the relationship into the momentum equation, one gets

$$\frac{d_r\vec{v}_r}{dt} = -2\vec{\Omega}\times\vec{v}_r - \frac{1}{\rho}\nabla p + \vec{g} + \Omega^2\vec{R}$$

Hereafter the subscript r is omitted. The effective gravity acceleration is defined as the sum of $\vec{g}$ and $\Omega^2\vec{R}$. Then the momentum equation in the rotating frame is

$$\frac{\partial\vec{v}}{\partial t} = -\vec{v}\cdot\nabla\vec{v} - 2\vec{\Omega}\times\vec{v} - \frac{1}{\rho}\nabla p + \vec{g}$$

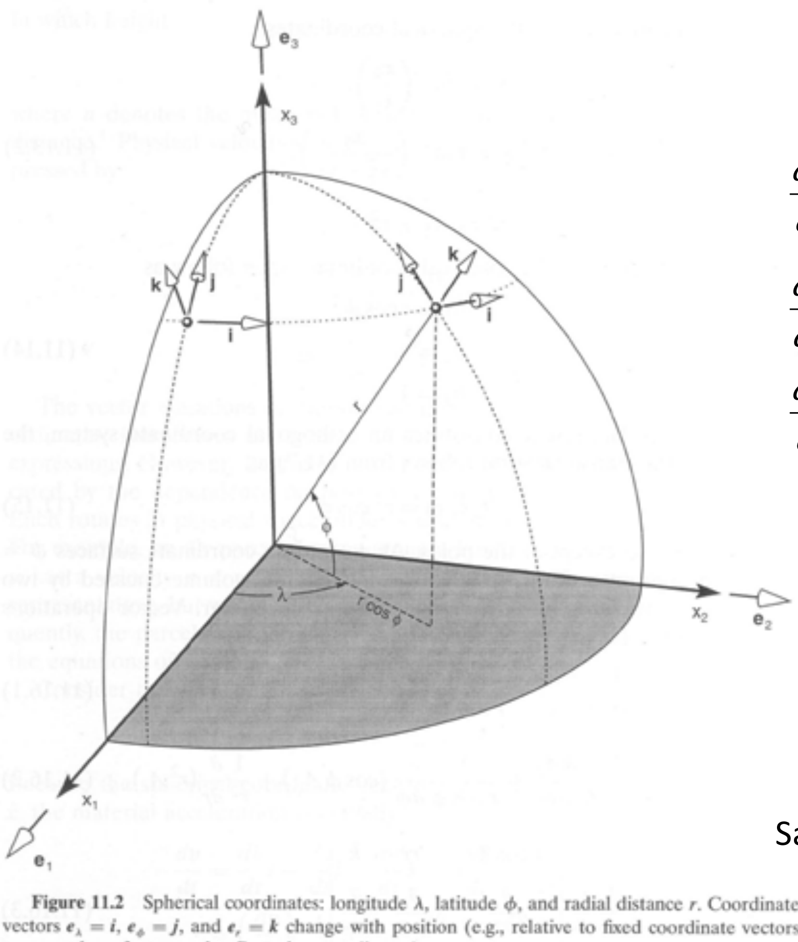
Next we consider an expression in a Cartesian coordinate system on a spherical surface. Using the velocity component (u, v, w) in the unit vector system $(\vec{i}, \vec{j}, \vec{k})$,

$$\vec{v} = \vec{i}u + \vec{j}v + \vec{k}w$$

$\vec{i}$: eastward, $\vec{j}$: northward, $\vec{k}$: upward

Then

$$\frac{d\vec{v}}{dt} = \vec{i}\frac{du}{dt} + \vec{j}\frac{dv}{dt} + \vec{k}\frac{dw}{dt} + u\frac{d\vec{i}}{dt} + v\frac{d\vec{j}}{dt} + w\frac{d\vec{k}}{dt}$$



$$\frac{d\vec{i}}{dt} = u\left(\frac{\tan\phi}{r}\vec{j} - \frac{\vec{k}}{r}\right)$$

$$\frac{d\vec{j}}{dt} = -u\frac{\tan\phi}{r}\vec{i} - \frac{v}{r}\vec{k}$$

$$\frac{d\vec{k}}{dt} = \frac{u}{r}\vec{i} + \frac{v}{r}\vec{j}$$

Salby (1996)

Figure 11.2 Spherical coordinates: longitude λ, latitude φ, and radial distance r. Coordinate vectors $e_\lambda = i$, $e_\phi = j$, and $e_r = k$ change with position (e.g., relative to fixed coordinate vectors e_1, e_2, and e_3 of rectangular Cartesian coordinates).

Momentum equations (a: planetary radius)

$$\frac{du}{dt} = \left(2\Omega\sin\phi + \frac{u\tan\phi}{a}\right)v - \frac{uw}{a} - 2\Omega w\cos\phi - \frac{1}{\rho}\frac{\partial p}{\partial x}$$

$$\frac{dv}{dt} = -\left(2\Omega\sin\phi + \frac{u\tan\phi}{a}\right)u - \frac{vw}{a} - \frac{1}{\rho}\frac{\partial p}{\partial y}$$

$$\frac{dw}{dt} = \frac{u^2+v^2}{a} - 2\Omega u\cos\phi - \frac{1}{\rho}\frac{\partial p}{\partial z} - g$$

$f = 2\Omega\sin\phi$: Coriolis parameter

The shaded terms are "metric" terms arising from the spherical geometry. These are small on the Earth. (Venus is exception)

When dw/dt = 0 , the equation in the vertical direction reduces to the hydrostatic equilibrium:

$$\frac{\partial p}{\partial z} = -\rho g$$

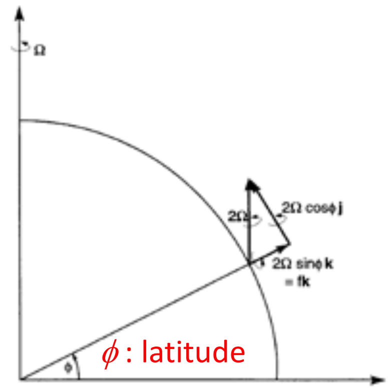


Figure 11.3 Planetary vorticity 2Ω decomposed into horizontal and vertical components.

Neglecting the metric terms and defining $\vec{v} = (u, v)$ as the horizontal component of the velocity, we have a set of governing equations:

$$\frac{\partial \vec{v}}{\partial t} = -\vec{v}\cdot\nabla\vec{v} - w\frac{\partial \vec{v}}{\partial z} - f\vec{k}\times\vec{v} - \frac{1}{\rho}\nabla p$$ horizontal momentum eq.

$$\frac{\partial w}{\partial t} = -\vec{v}\cdot\nabla w - w\frac{\partial w}{\partial z} - \frac{1}{\rho}\frac{\partial p}{\partial z} - g$$ vertical momentum eq.

$$\frac{\partial \rho}{\partial t} = -\nabla\cdot(\rho\vec{v}) - \frac{\partial(\rho w)}{\partial z}$$ continuity eq.

$$p = \rho R T$$ state eq.

$$\frac{\partial \theta}{\partial t} = -\vec{v}\cdot\nabla\theta - w\frac{\partial \theta}{\partial z} + \underbrace{\frac{1}{c_p}\left(\frac{p_S}{p}\right)^{R/C_p}\dot{H}}_{\text{diabatic heating}}$$ thermodynamics eq.

$\vec{k}$: unit vector in the vertical direction

$f = 2\Omega\sin\phi$: Coriolis parameter

$\theta \equiv T\left(\frac{p_S}{p}\right)^{R/C_p}$: potential temperature

$\dot{H}$: diabatic heating rate ($J/m^3/s$)

Pressure coordinate

Large-scale atmospheric motions satisfy hydrostatic equilibrium.
In this case the pressure p can be used as the vertical coordinate:

$$p = p(z) \quad \rightarrow \quad \boxed{z = z(p)}$$

p : pressure, z : altitude

$\boxed{\omega} \equiv dp/dt$: vertical velocity in pressure coordinate

$\boxed{\Phi}$: geopotential $\quad d\Phi = g dz$

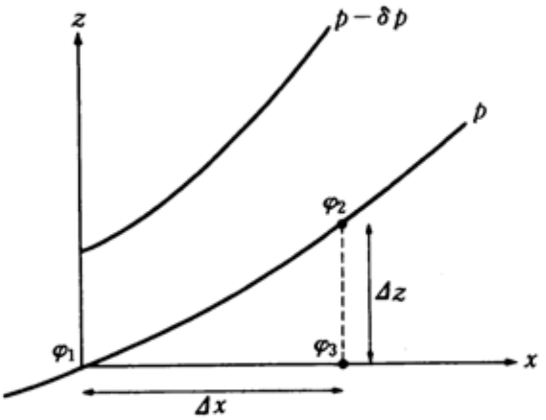


図 3.2 x–z 面内の断面. 小倉(1978)

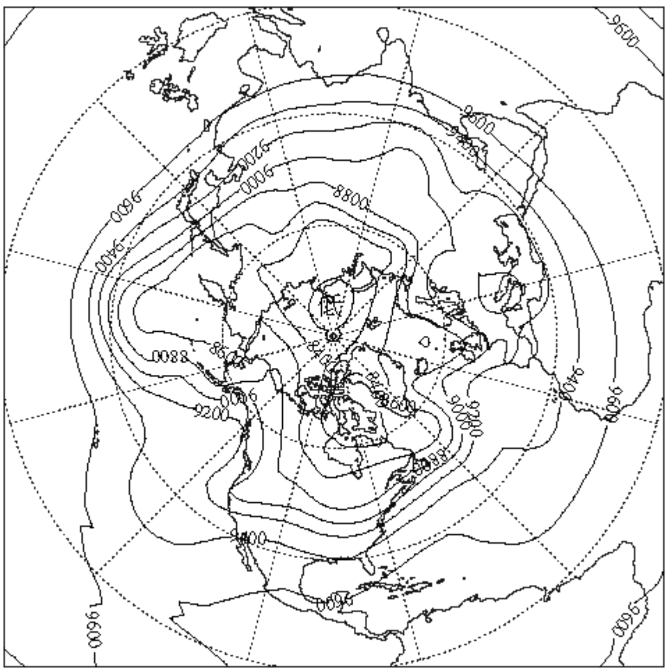
Approximate relationship between ω and w:

$$\omega \equiv dp/dt$$
$$\sim -w/\rho g$$

Pressure gradient force in the momentum equation:

$$-\frac{1}{\rho}\nabla p \quad \Rightarrow \quad -\nabla\Phi$$

300hPa, 2002/1/15



Primitive equations:

$$\frac{\partial \vec{v}}{\partial t} = -\vec{v}\cdot\nabla\vec{v} - \omega\frac{\partial \vec{v}}{\partial p} - f\vec{k}\times\vec{v} - \nabla\Phi$$ horizontal momentum eq.

$$\frac{\partial \Phi}{\partial p} = -\frac{RT}{p}$$ state eq. + hydrostatic equilibrium

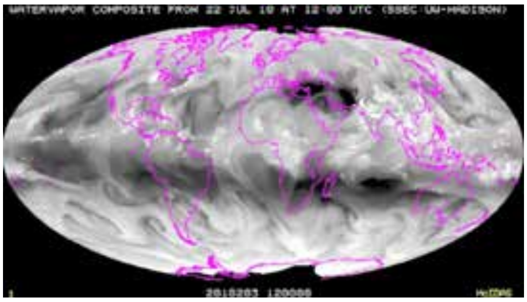
$$\nabla\cdot\vec{v} + \frac{\partial \omega}{\partial p} = 0$$ continuity eq.

$$\frac{\partial \theta}{\partial t} = -\vec{v}\cdot\nabla\theta - \omega\frac{\partial \theta}{\partial p} + \frac{1}{c_p}\left(\frac{p_S}{p}\right)^{R/C_p}\dot{H}$$ thermodynamics eq.

$\vec{v} = (u, v)$: horizontal velocity
ω : vertical velocity
$\vec{k}$: unit vector in the vertical direction
$f = 2\Omega\sin\phi$: Coriolis parameter
$\dot{H}$: diabatic heating rate (J/m^3/s)

Geostrophic flow

For synoptic scale (>1000 km) motions in Earth's atmosphere:
horizontal scale L ~ 1000 km
pressure scale P ~ 1000 hPa
velocity scale U ~ 10 m/s
time scale L/U ~ 10^5 s
Coriolis parameter f ~ 10^{-4} Hz



$$\frac{\partial \vec{v}}{\partial t} = -\vec{v}\cdot\nabla\vec{v} - \omega\frac{\partial \vec{v}}{\partial p} - f\vec{k}\times\vec{v} - \nabla\Phi$$

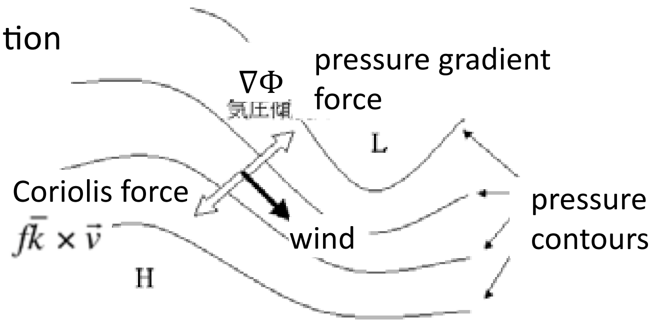
Coriolis force; pressure gradient force

10^{-4} 10^{-4} 10^{-4} 10^{-3} 10^{-3} (m s^{-2})

The ratio between the acceleration (U^2/L) and the Coriolis force (fU)

$$R_o \equiv \frac{U}{fL}$$: Rossby number (around 0.1 on Earth, 10–100 on Venus)

Geostrophic flow is a good approximation for small Rossby numbers.



Definition of geostrophic flow

$$f\bar{k} \times \vec{v}_g = -\nabla\Phi$$

$$\vec{v}_g = \frac{1}{f}\vec{k} \times \nabla\Phi \qquad \text{or} \qquad u_g = -\frac{1}{f}\frac{\partial\Phi}{\partial y}, \quad v_g = \frac{1}{f}\frac{\partial\Phi}{\partial x}$$

Geostrophic flow is two-dimensional:

$$\nabla\cdot\vec{v}_g = 0$$

$$\omega = 0$$

Geostrophic flow is not valid when Ro (= U/fL) is large, for example:

- near the equator (f is small)
- scales < O(100km) (L is small)
- winds are fast like Mars (U is large)
- planetary rotation is slow like Venus (f is small)

Thermal wind

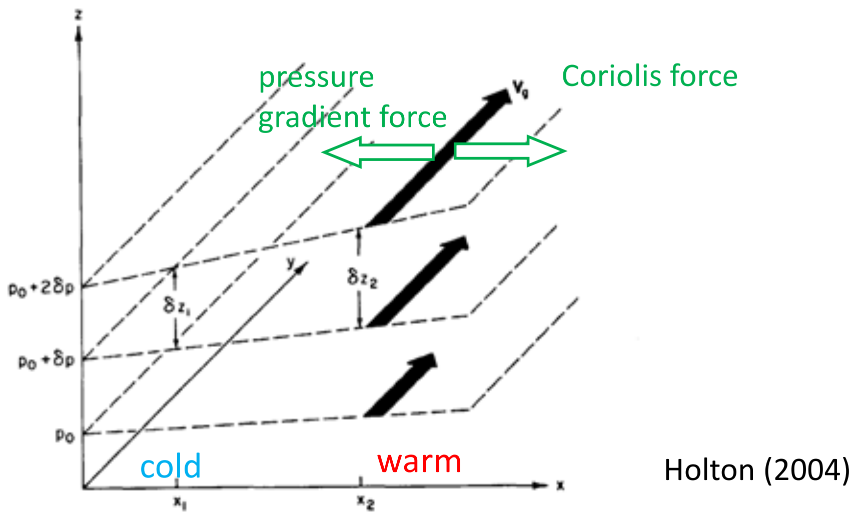
Differentiating $f\bar{k} \times \vec{v}_g = -\nabla\Phi$ with respect to p and using $\frac{\partial\Phi}{\partial p} = -\frac{RT}{p}$, "thermal wind" relation is obtained:

$$\frac{\partial\vec{v}_g}{\partial p} = -\frac{R}{fp}\vec{k} \times \nabla T$$

or

$$\frac{\partial u_g}{\partial p} = \frac{R}{fp}\frac{\partial T}{\partial y}$$

$$\frac{\partial v_g}{\partial p} = -\frac{R}{fp}\frac{\partial T}{\partial x}$$

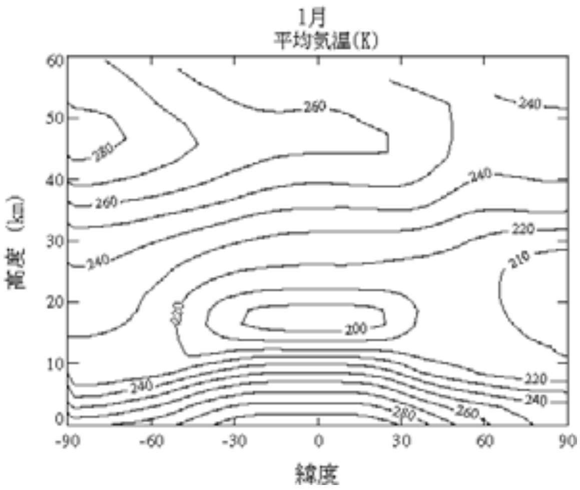


Holton (2004)

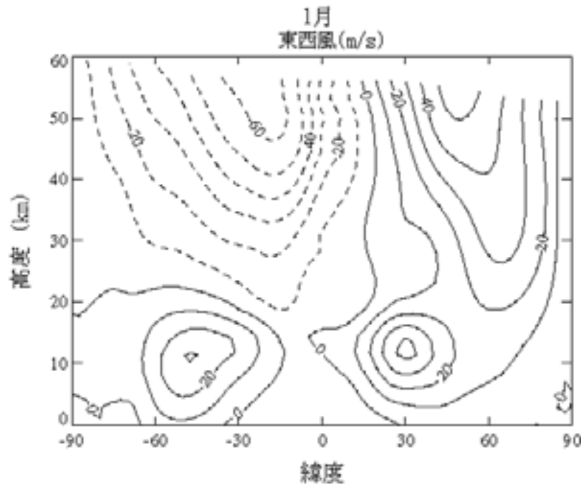
Fig. 3.8 Relationship between vertical shear of the geostrophic wind and horizontal temperature gradients. (Note: $\delta p < 0$.)

Latitude-altitude cross section of Earth's atmosphere

Temperature



Eastward wind



Latitude-altitude cross section of Mars's atmosphere

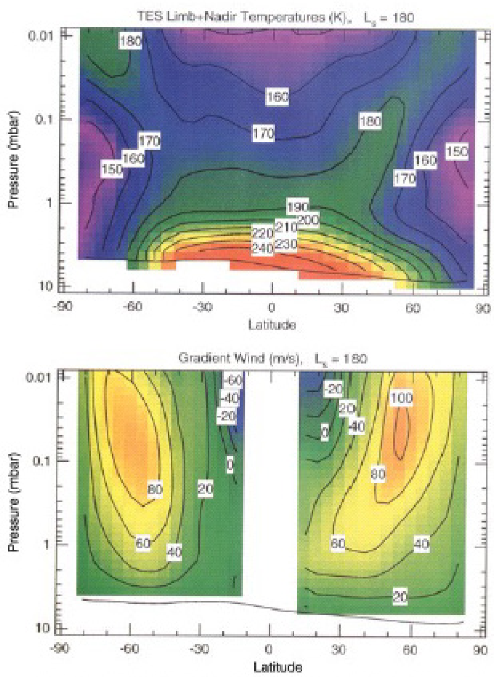


Plate 6. Mean meridional cross sections of (top) retrieved atmospheric temperatures and (bottom) gradient wind speeds for $L_s = 180°$ (northern hemisphere fall equinox). Temperatures are in kelvins, and wind speeds are in m s^{-1}. Positive wind speeds are eastward. The mean surface pressures for this data set are indicated at the lower boundary of the contoured regimes.

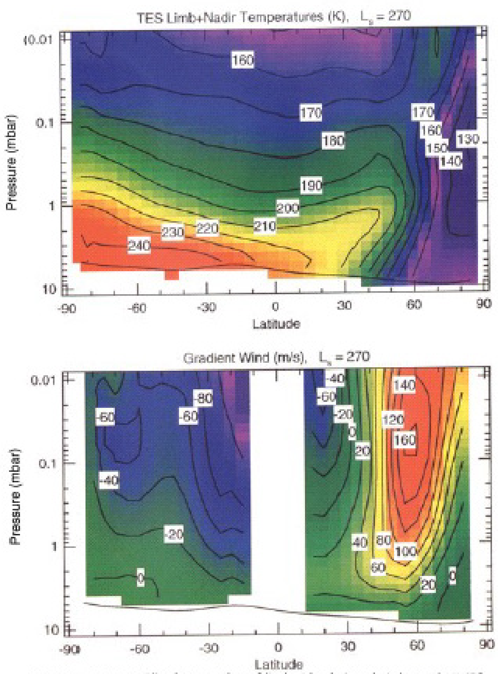


Plate 7. Mean meridional cross sections of (top) retrieved atmospheric temperatures and (bottom) gradient wind speeds for $L_s = 270°$ (northern winter solstice). Temperatures are in kelvins, and wind speeds are in m s^{-1}. Positive wind speeds are eastward. The mean surface pressures for this data set are indicated at the lower boundary of the contoured regimes.

Temperature

Eastward wind

(Smith et al. 2001)

Geostrophic flow on Mars?

For synoptic scale (>1000 km) motions in Earth's atmosphere:
- horizontal scale L ~ 1000 km
- pressure scale P ~ 1000 hPa
- velocity scale U ~ 10 m/s on Earth → 100 m/s on Mars (westerly jet)
- time scale L/U ~ 10^5 s → 10^4 s on Mars
- Coriolis parameter f ~ 10^{-4} Hz

$$\cancel{\frac{\partial \vec{v}}{\partial t}} = \cancel{-\vec{v}\cdot\nabla\vec{v}} \cancel{- \omega\frac{\partial \vec{v}}{\partial p}} - \underline{f\vec{k}\times\vec{v}} - \underline{\nabla\Phi}$$

Coriolis force; pressure gradient force

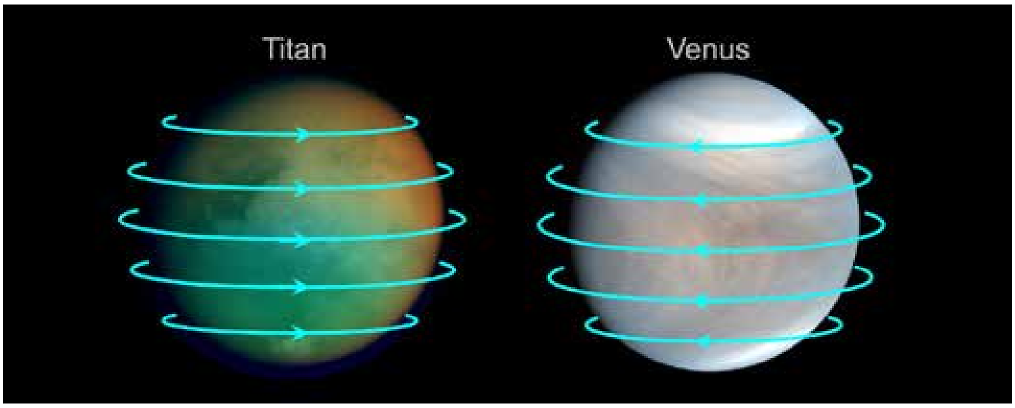
10^{-4}	10^{-4}	10^{-4}	10^{-3}	10^{-3}	(m s^{-2})	on Earth
10^{-2}	10^{-2}	10^{-2}	10^{-2}	10^{-2}	(m s^{-2})	on Mars

The ratio between the acceleration (U^2/L) and the Coriolis force (fU)

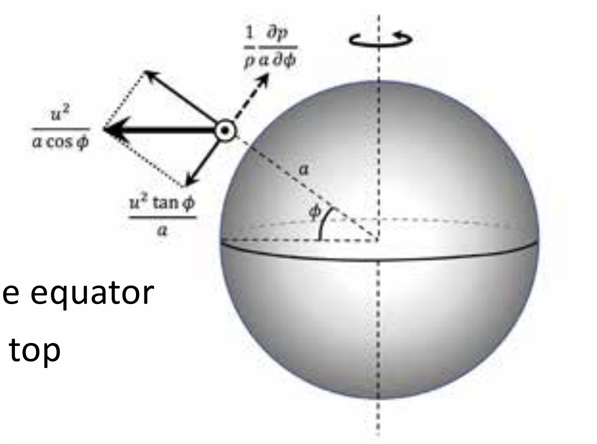
$$R_o \equiv \frac{U}{fL}$$

: Rossby number (around 0.1 on Earth, 10–100 on Venus)
around 1 on Mars ?

Superrotation



Superrotation



It prevails in Venus's atmosphere

planetary rotation: 1.8 m/s on the equator

zonal wind: 100 m/s at the cloud top

Momentum equations

$$\frac{du}{dt} = \left(2\Omega\sin\phi + \frac{u\tan\phi}{a}\right)v - \frac{uw}{a} - 2\Omega w\cos\phi - \frac{1}{\rho}\frac{\partial p}{\partial x}$$

$$\frac{dv}{dt} = -\left(2\Omega\sin\phi + \frac{u\tan\phi}{a}\right)u - \frac{vw}{a} - \frac{1}{\rho}\frac{\partial p}{\partial y}$$

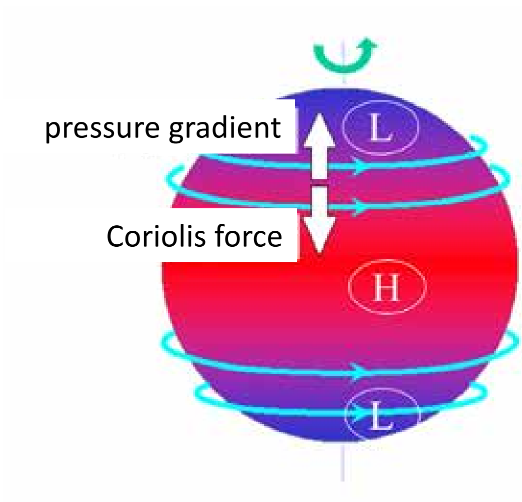
$$\frac{dw}{dt} = \frac{u^2+v^2}{a} - 2\Omega u\cos\phi - \frac{1}{\rho}\frac{\partial p}{\partial z} - g$$

Atmospheric rotation takes the place of planetary rotation.

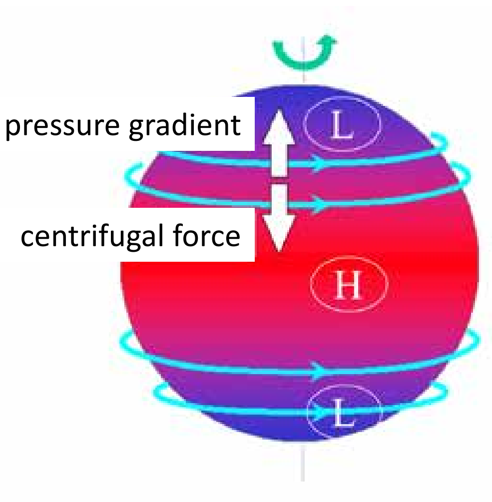
Meridional force balance of zonal flow

geostrophic flow (planetary rotation >> wind)

cyclostrophic flow (planetary rotation << wind)



Earth-like



Venus-like

Thermal wind

geostrophic flow
(planetary rotation >> wind)

cyclostrophic flow
(planetary rotation << wind)

$$\cancel{\frac{dv}{dt}} = -\left(\boxed{2\Omega\sin\phi} + \boxed{\frac{u\tan\phi}{a}}\right)u - \cancel{\frac{vw}{a}} - \frac{1}{\rho}\frac{\partial p}{\partial y}$$

$$2\Omega\sin\phi\; a\; u + \frac{1}{\rho}\frac{\partial p}{\partial \phi} = 0$$

$$u^2\tan\phi + \frac{1}{\rho}\frac{\partial p}{\partial \phi} = 0$$

← state eq.
← hydrostatic equilibrium

$$\frac{\partial u}{\partial \zeta} + \frac{R}{2a\Omega\sin\phi}\frac{\partial T}{\partial \phi} = 0$$

$$\frac{\partial u^2}{\partial \zeta} + \frac{R}{\tan\phi}\frac{\partial T}{\partial \phi} = 0$$

$\zeta = -\ln p$: log-pressure altitude

Latitude-altitude cross section of Venus's atmosphere

Temperature

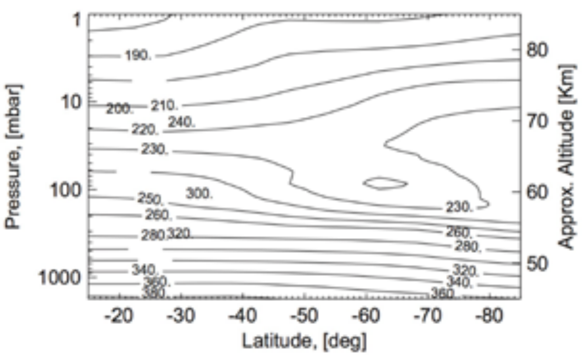


Fig. 2. Meridional cross section of atmospheric temperature (K) obtained combining 116 VeRa profiles. Hemispherical symmetry and local time independence has been assumed. Contours have been smoothed. Contours interval is 10 K, some contours have been removed to render the plot clearer.

Westward wind

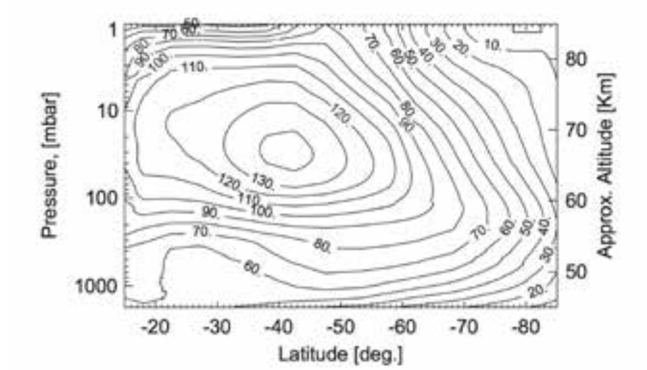
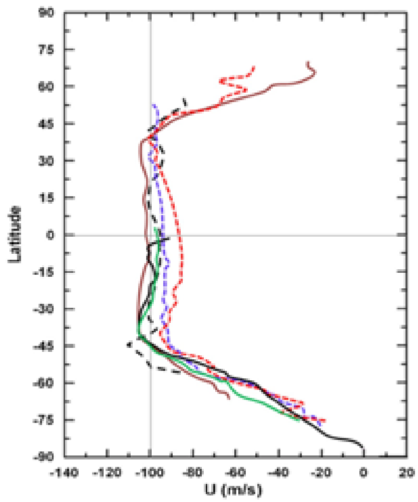


Fig. 8. Latitude–height cross section of zonal thermal wind speed ($m\,s^{-1}$) derived from VeRa temperature profiles assuming cyclostrophic balance (Eq. (5)). The velocity used as lower boundary condition is the cloud-tracked wind profile retrieved from VIRTIS/VEx 1.74 μm images (Sánchez-Lavega et al., 2008). Contour interval is 10 $m\,s^{-1}$.

(Piccialli et al., 2012)

Latitude-altitude cross section of Venus's atmosphere

Cloud-tracked zonal winds around 70 km (Sánchez-Lavega et al. 2017)



Zonal winds determined from cyclostrophic balance (Piccialli et al., 2012)

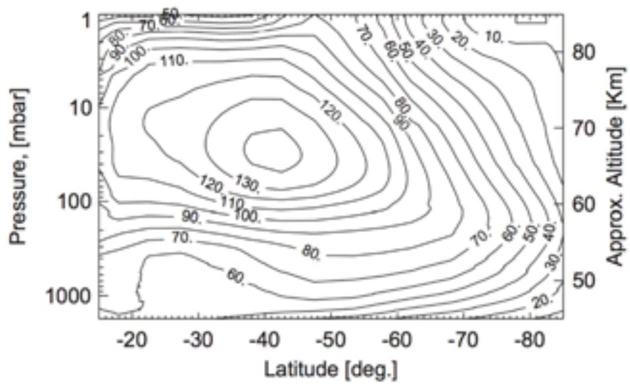


Fig. 8. Latitude–height cross section of zonal thermal wind speed (m s^{-1}) derived from VeRa temperature profiles assuming cyclostrophic balance (Eq. (5)). The velocity used as lower boundary condition is the cloud-tracked wind profile retrieved from VIRTIS/VEx 1.74 μm images (Sánchez-Lavega et al., 2008). Contour interval is 10 m s^{-1}.

The discrepancy is not understood

Radiative energy budget and meridional circulation

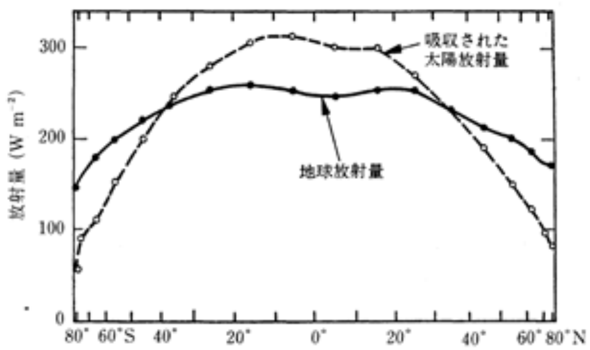


図 4. 4 1962～66 年の衛星観測に基づく帯状平均した放射収支の緯度分布（Vonder Haar and V. Suomi, 1971：*J. Atmos. Sci.*, 28, 305-314.）

小倉(1978)

Hadley cell

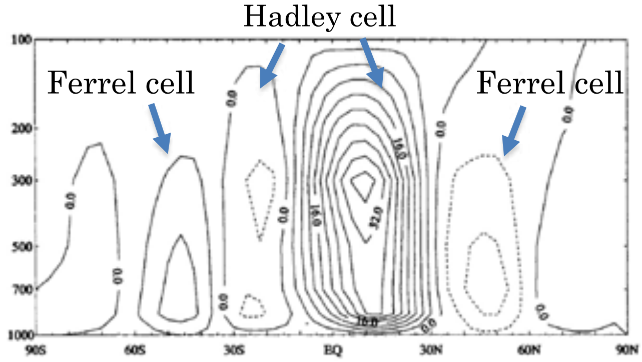


Fig. 10.7 Streamfunction (units: 10^2 kg m^{-1}s^{-1}) for the observed Eulerian mean meridional circulation for Northern Hemisphere winter, based on the data of Schubert et al. (1990).

Holton (2004)

Horizontal convection →Hadley circulation

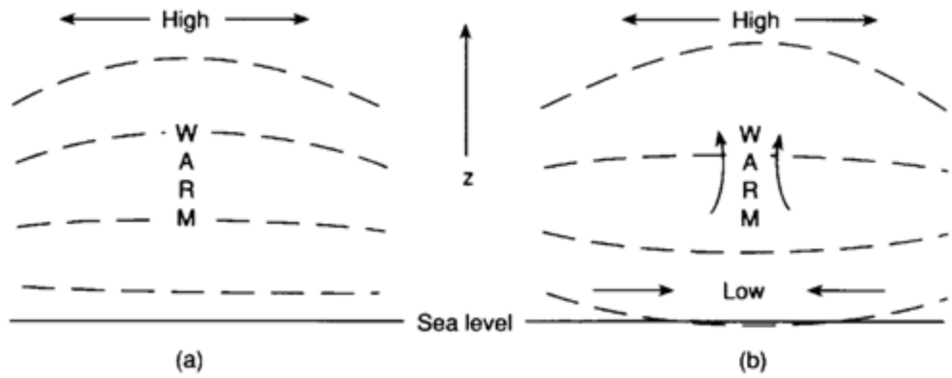
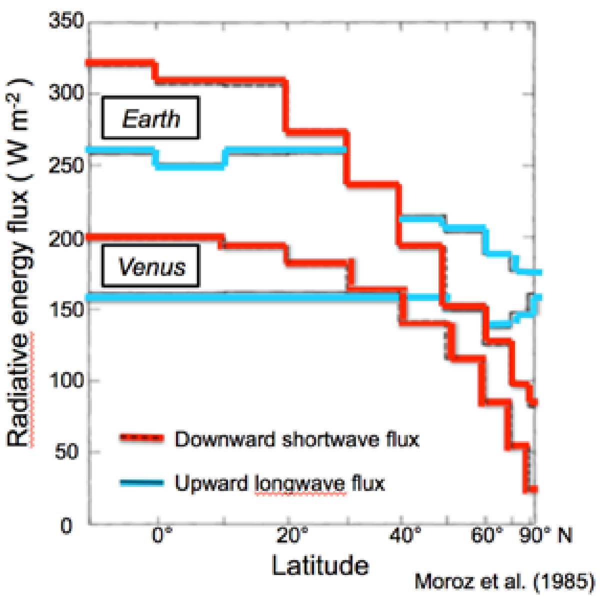
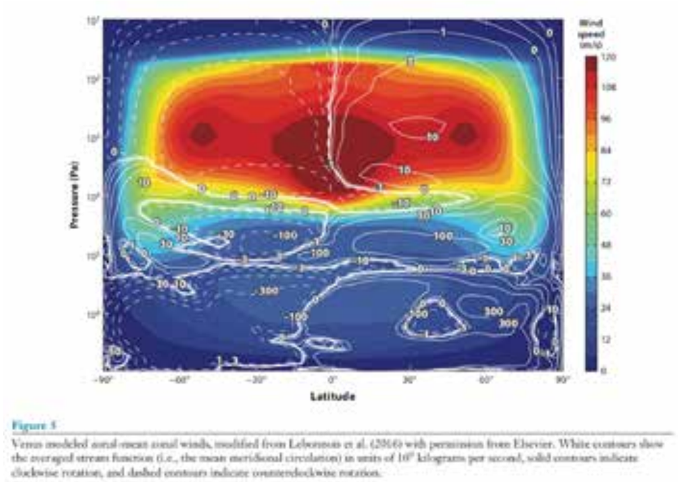


Fig. 3.11 Adjustment of surface pressure to midtropospheric heat source. Dashed lines indicate isobars.

Holton (1992)

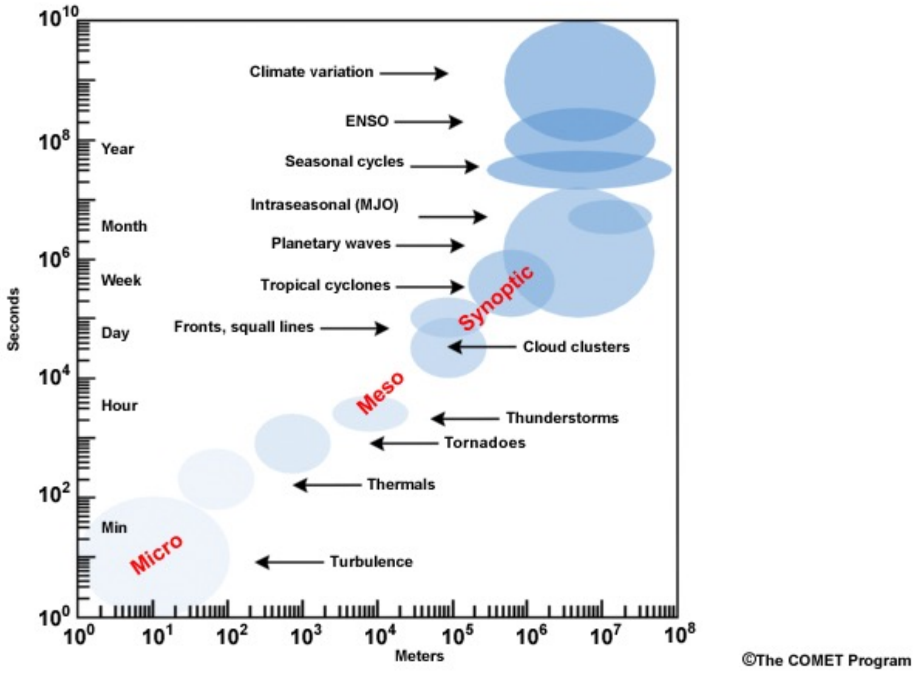
Radiative energy budget and meridional circulation





Read & Lebonnois (2018)

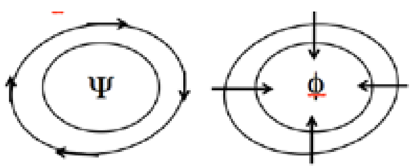
Space and time-scales of dynamical atmospheric processes



Rotational wind and divergent wind

The horizontal velocity vector can be expressed with the stream function ψ and the velocity potential ϕ :

$$\vec{v} = \underbrace{\vec{k} \times \nabla\psi}_{\text{Rotational wind}} + \underbrace{\nabla\phi}_{\text{Divergent wind}}$$

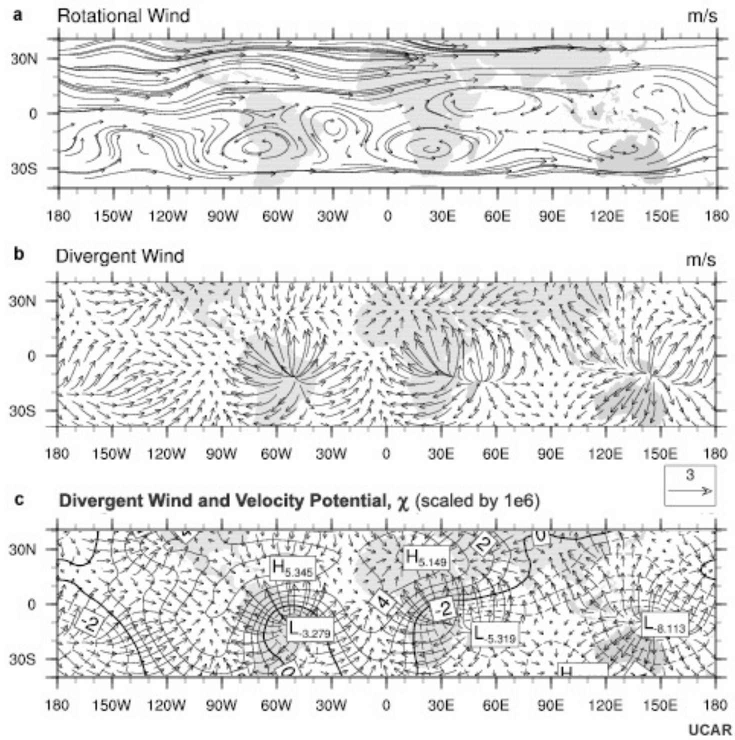


Rotational component is nondivergent

$$\nabla\cdot(\vec{k} \times \nabla\psi) = 0$$: geostrophic flow, Rossby wave

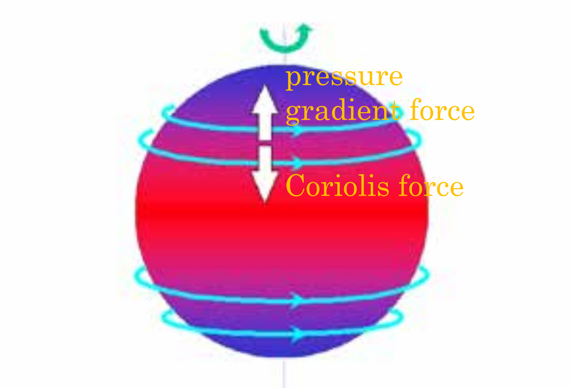
Divergent component is irrotational

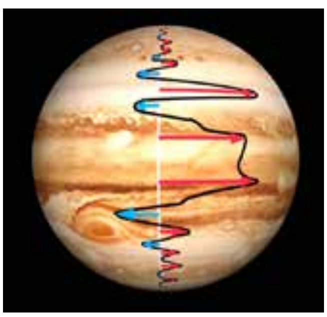
$$\nabla \times \nabla\phi = 0$$: convection, gravity wave



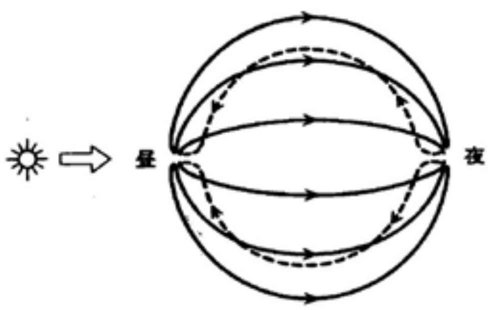
Planetary-scale motions

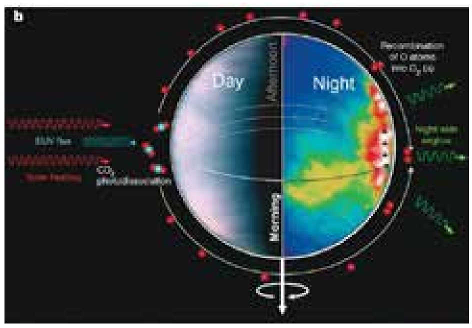
(rotational flow)
Mean zonal wind





(divergent flow)
Hadley circulation
Subsolar-to-antisolar circulation

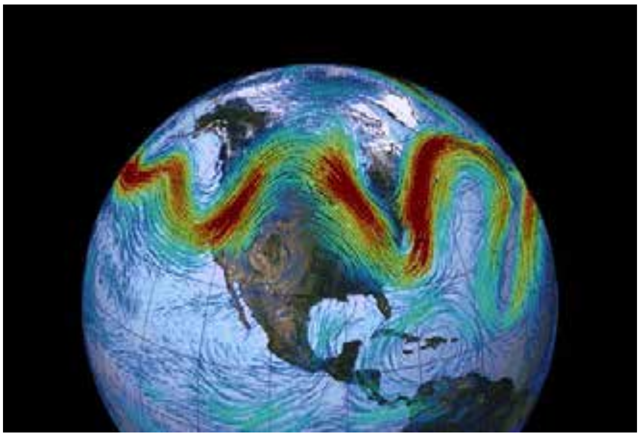




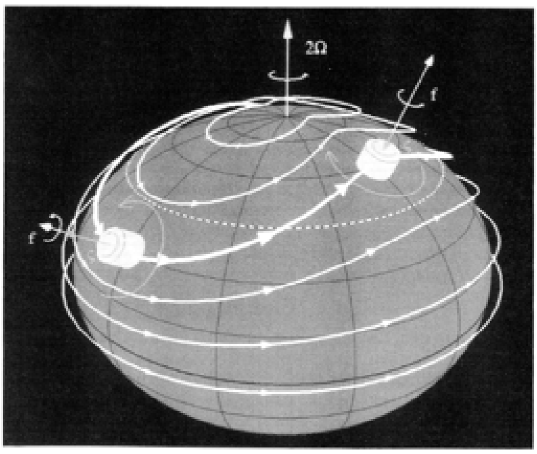
Atmospheric waves

- generated in an unstable background atmosphere
- can transport momentum and energy over long distances
- can induce mixing of the atmosphere

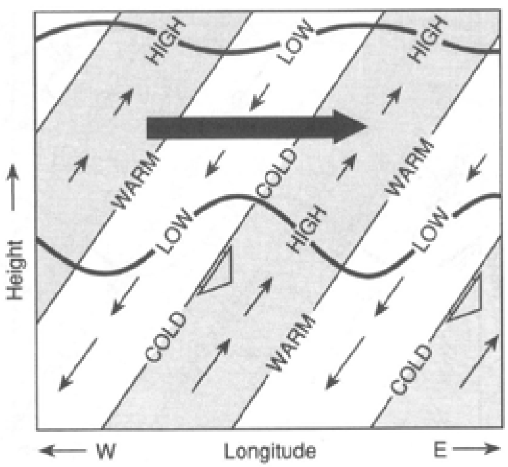
→ Waves play crucial roles in the development of planetary-scale atmospheric structure



Rossby wave (horizontal oscillation)



Gravity wave（vertical oscillation）



Salby (1996)

Vorticity

The *circulation C* about a closed contour is defined as the line integral along the contour of the component of the velocity vector that is locally tangent to the contour:

$$C = \oint_l \vec{v} \cdot d\vec{l} = \oint_l |\vec{v}| \cos\alpha dl = \iint_S (\nabla \times \vec{v}) \cdot \vec{n} dS$$

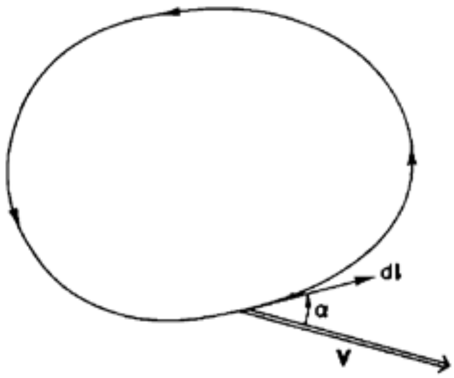
n : unit vector normal to the surface

(Stokes' theorem was used.)

Dividing the circulation δC encircling a small area δS, and considering the limit $\delta S \rightarrow 0$, we get the *vorticity*:

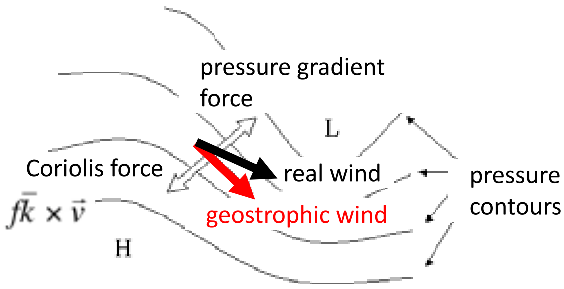
$$\zeta = \frac{\delta C}{\delta S} = (\nabla \times \vec{v}) \cdot \vec{k}$$

$$= \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y}$$



Holton (1992)

Quasi-geostrophic approximation



The real wind is divided into geostrophic wind and ageostrophic wind.

$$\vec{v} = \boxed{\vec{v}_g} + \boxed{\vec{v}_a}$$

$$\vec{v}_g \equiv \frac{1}{f_0} \vec{k} \times \nabla \Phi$$

$$\frac{v_a}{v_g} = O(R_o) \ll 1$$

$$R_o \equiv \frac{U}{fL}$$

The ratio of the magnitudes of the ageostrophic and geostrophic winds is the same order of magnitude as the Rossby number (around 0.1 in Earth's atmosphere)

Lagrangian derivative: $$\frac{d}{dt} = \frac{\partial}{\partial t} + \vec{v} \cdot \nabla + \omega \frac{\partial}{\partial p}$$

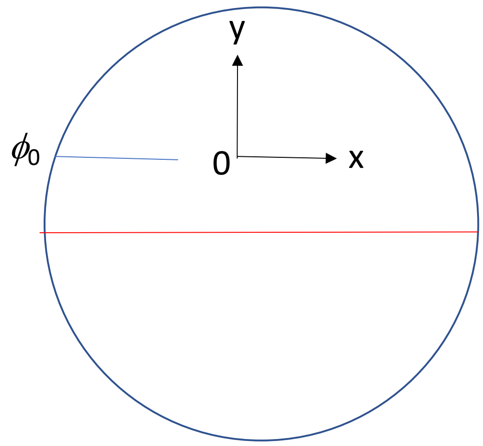
The advection term is approximated by advection by geostrophic wind:

$$\frac{d_g}{dt} \equiv \frac{\partial}{\partial t} + \boxed{\vec{v}_g} \cdot \nabla$$

Beta-plane approximation

The first-order Taylor series approximation of the Coriolis parameter f :

$$
\begin{aligned}
f &= 2\Omega \sin\phi \\
&\sim 2\Omega\,[\sin\phi_0 + (\phi - \phi_0)\cos\phi_0] \\
&= \boxed{f_0 + \beta y}
\end{aligned}
$$



where

$$f_0 = 2\Omega \sin\phi_0$$

$$\beta = \frac{2\Omega}{a}\cos\phi_0 \sim \frac{df}{dy}$$

Mid-latitude beta-plane : $f = f_0 + \beta y$

Equatorial beta-plane : $f = \beta y$

Quasi-geostrophic vorticity equation

Introducing beta-plane approximation to the quasi-geostrophic Lagrangian derivative and retaining small quantities to the first order, the rate of change of the geostrophic wind is given by

$$
\begin{aligned}
\left(\frac{\partial}{\partial t} + \vec{v}_g \cdot \nabla\right)\vec{v}_g &= -(f_0 + \beta y)\vec{k} \times (\vec{v}_g + \vec{v}_a) - \nabla\Phi \\
&\sim -f_0\vec{k} \times \vec{v}_a - \beta y\vec{k} \times \vec{v}_g \qquad (1)
\end{aligned}
$$

($\vec{v}_g \equiv \frac{1}{f_0}\vec{k} \times \nabla\Phi$ was used)

Since geostrophic wind is nondivergent ($\nabla \cdot \vec{v}_g = 0$), the continuity eq. is

$$\nabla \cdot \vec{v}_a + \frac{\partial \omega}{\partial p} = 0 \qquad (2)$$

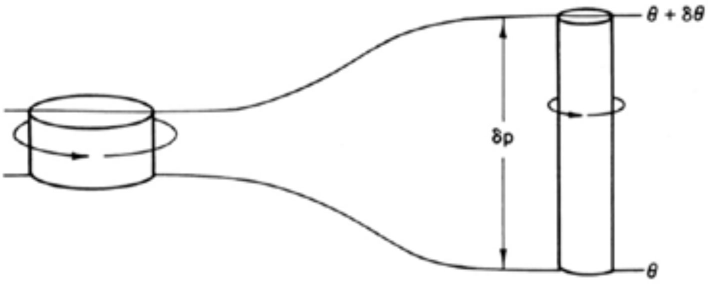
Operating rotation (∇×) to (1) and using (2), we have the quasi-geostrophic vorticity equation.

Quasi-geostrophic vorticity equation:

$$\frac{\partial \zeta_g}{\partial t} = -\vec{v}_g \cdot \nabla(\zeta_g + f) + f_0 \frac{\partial \omega}{\partial p}$$

$$\zeta_g \equiv \frac{\partial v_g}{\partial x} - \frac{\partial u_g}{\partial y} = \frac{\nabla^2 \Phi'}{f_0}$$: geostrophic vorticity

Vorticity changes with time through
- advection of absolute vorticity ($\varsigma_g + f$) by geostrophic wind ($\vec{v}_g$)
- vertical divergence (horizontal divergence)



Holton (2004)

Fig. 4.7 A cylindrical column of air moving adiabatically, conserving potential vorticity.

Rossby wave

Let us consider a two-dimensional motion (ω = 0)

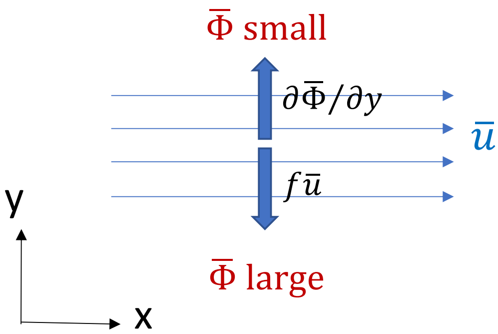
$$\left(\frac{\partial}{\partial t} + \vec{v}_g \cdot \nabla\right)(\zeta_g + f) = 0$$

→ Absolute vorticity ($\varsigma_g + f$) is conserved along the geostrophic wind $\vec{v}_g$.

A basic state where a homogeneous zonal flow exists:

$$f\bar{u} = -\partial\bar{\Phi}/\partial y$$

$$\bar{\zeta}_g = 0$$



The deviation from the basic state is denoted by ()' :

$$\frac{\partial}{\partial t}\zeta_g + (\bar{u} + u'_g)\frac{\partial \zeta_g}{\partial x} + v'_g \frac{\partial \zeta_g}{\partial y} + v'_g \beta = 0 \qquad (1)$$

The velocity and vorticity are related to the geopotential perturbation

$$\zeta_g = \frac{\nabla^2 \Phi'}{f_0}, \; u'_g = -\frac{1}{f_0}\frac{\partial \Phi'}{\partial y}, \; v'_g = \frac{1}{f_0}\frac{\partial \Phi'}{\partial x}$$

Substituting these into (1) and retaining first order terms only, we get

$$\frac{\partial}{\partial t}\nabla^2\Phi' + \bar{u}\frac{\partial}{\partial x}\nabla^2\Phi' + \beta\frac{\partial \Phi'}{\partial x} = 0$$

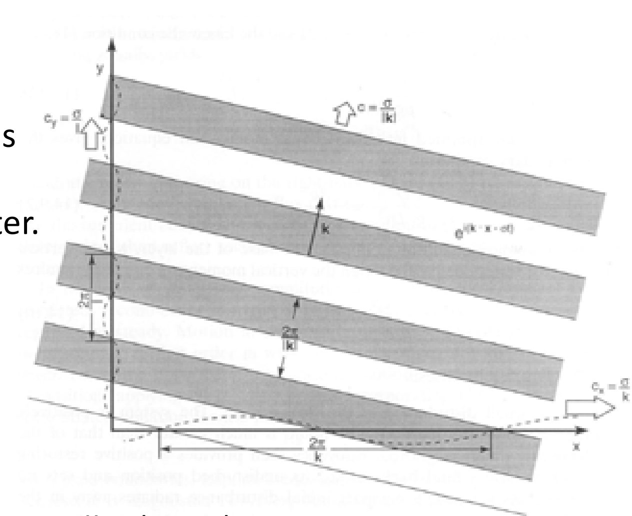
Assuming a wave solution $\Phi' = \hat{\Phi}\exp[i(kx + ly - kct)]$, the phase velocity of Rossby wave is obtained as:

$$c - \bar{u} = -\frac{\beta}{k^2 + l^2}$$

k: zonal wavenumber
l: meridional wavenumber
c: zonal phase velocity

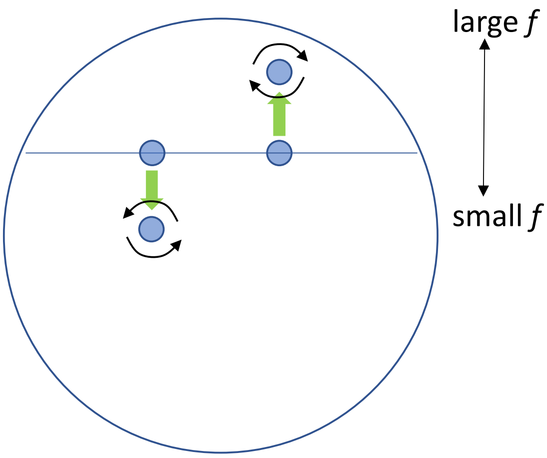
- Propagation opposite to the planetary rotation
- β effect (latitude variation of the Coriolis parameter f) is needed.
- Longer waves (smaller k) propagate faster.

The wave possesses angular momentum in the direction opposite to the planetary rotation



Salby (1996)

Propagation of Rossby waves



$$\left(\frac{\partial}{\partial t} + \vec{v}_g \cdot \nabla\right)(\zeta_g + f) = 0$$

$$f = 2\Omega \sin\phi$$

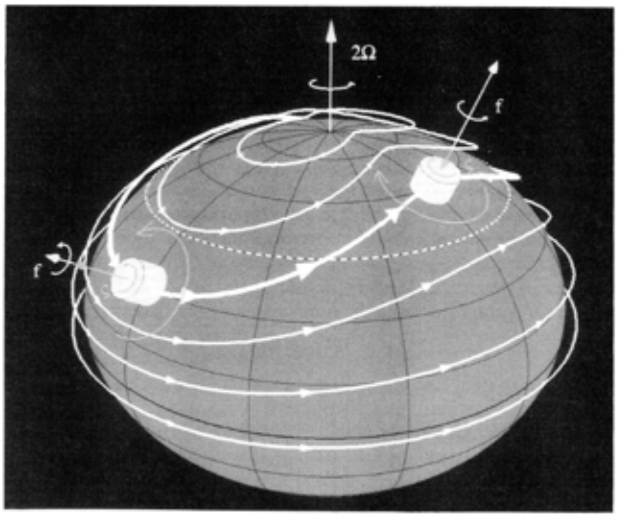
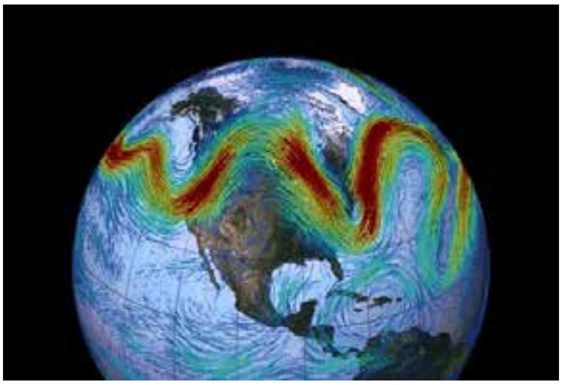


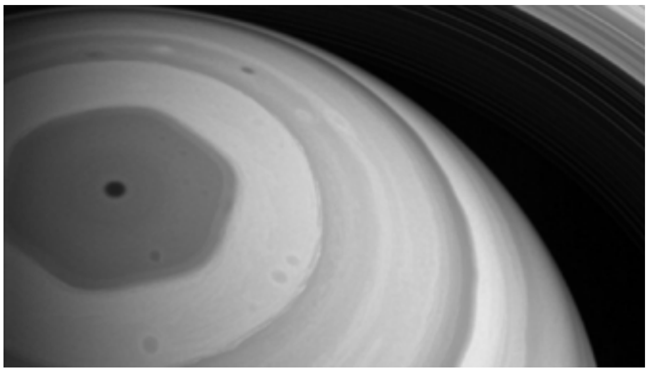
Figure 14.16 Schematic illustrating the reaction of an air parcel to meridional displacement. Displaced equatorward, an eastward-moving parcel spins up cyclonically to conserve absolute vorticity. Northward motion induced ahead of it then deflects the parcel's trajectory poleward back toward its undisturbed latitude. The reverse process occurs when the parcel overshoots and is displaced poleward of its undisturbed latitude.

Salby (1996)

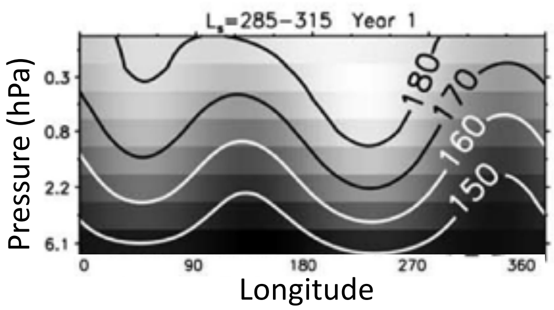
Rossby waves in planetary atmospheres



Earth



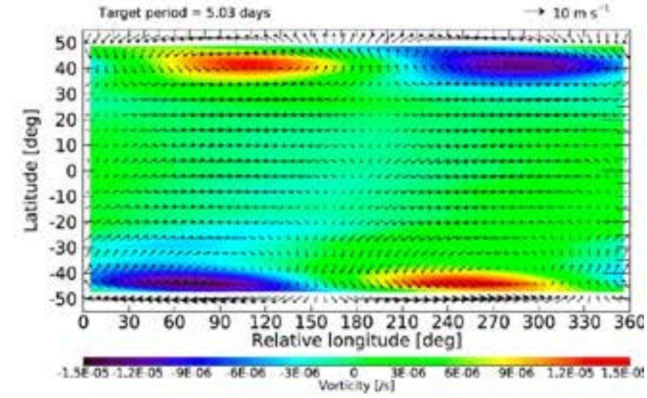
Saturn

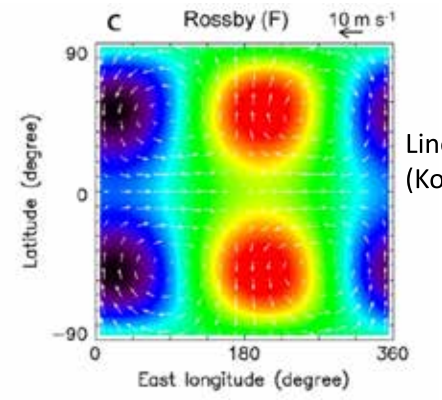


Mars (MGS/TES temperature)

Rossby waves on Venus

Venus' 5-day wave observed by cloud-tracking
(Imai et al. 2019)





Linear solution of Rossby wave at 70 km
(Kouyama et al. 2015)

The superrotation of the atmosphere takes the place of planetary rotation.

Gravity wave

In Cartesian coordinates, without the assumption of hydrostatic equilibrium, the governing equations are:

$$\frac{\partial u}{\partial t} + u\frac{\partial u}{\partial x} + w\frac{\partial u}{\partial z} = \frac{1}{\rho}\frac{\partial p}{\partial x}$$

horizontal momentum eq.
(x-axis only)

$$\frac{\partial w}{\partial t} + u\frac{\partial w}{\partial x} + w\frac{\partial w}{\partial z} = -\frac{1}{\rho}\frac{\partial p}{\partial z}$$

vertical momentum eq.

$$\frac{\partial \rho}{\partial t} = -\frac{\partial(\rho u)}{\partial x} - \frac{\partial(\rho w)}{\partial z}$$

continuity eq.

$$\frac{\partial \theta}{\partial t} + u\frac{\partial \theta}{\partial x} + w\frac{\partial \theta}{\partial z} = 0$$

thermodynamics eq.

$$\theta = \underbrace{\boxed{\frac{p}{\rho R}}}_{T}\left(\frac{p_S}{p}\right)^{R/C_p}$$

Equations for disturbances:

$$\bar{\rho}\frac{\partial u'}{\partial t} = -\frac{\partial p'}{\partial x}$$

horizontal momentum eq. (x-axis only)

$$\bar{\rho}\frac{\partial w'}{\partial t} = -\frac{\partial p'}{\partial z} - \rho' g$$

vertical momentum eq.

$$\frac{\partial \rho'}{\partial t} = -\bar{\rho}\frac{\partial u'}{\partial x} - \frac{\partial(\bar{\rho} w')}{\partial z}$$

continuity eq.

$$\frac{1}{\bar{\theta}}\frac{\partial \theta'}{\partial t} + w'\frac{N^2}{g} = 0$$

thermodynamics eq.

$$\rho' = \frac{p'}{c_s^2} - \bar{\rho}\frac{\theta'}{\bar{\theta}}$$

$$c_s^2 = \frac{c_p}{c_v}RT$$

c_s : sound speed

$$N^2 = g\frac{d\ln\bar{\theta}}{dz}$$

N : buoyancy frequency

Assuming an isothermal atmosphere:

$$N^2 = g/c_p T$$

$$\bar{\rho}(z) = \rho_s \exp(-z/H)$$

Substituting the wave solution $w'(x,z,t) = \hat{w}(z)\exp[i(kx+\sigma t)]$ (σ : frequency) into the governing equations before, an equation for the vertical velocity *w* is obtained as

$$\frac{d^2(\bar{\rho}\hat{w})}{dz^2} + \frac{1}{H}\frac{d(\bar{\rho}\hat{w})}{dz} + \left[\frac{\sigma^2}{c_s^2} - k^2 + \frac{N^2k^2}{\sigma^2}\right]\bar{\rho}\hat{w} = 0$$

Considering the amplitude growth with height in a stratified atmosphere, w is assumed to have the form

$$\hat{w}(z) = W(z)\exp(z/2H) \quad \leftrightarrow \quad \bar{\rho}\hat{w}^2(z) \propto W^2(z)$$

Then the equation becomes

$$\frac{d^2W}{dz^2}+\left[\frac{\sigma^2}{c_s^2}-k^2+\frac{N^2k^2}{\sigma^2}-\frac{1}{4H^2}\right]W=0$$

Assuming a wave solution $W \propto \exp(imz)$ (*m*: vertical wavenumber), the dispersion relation is obtained:

$$m^2=\frac{\sigma^2}{c_s^2}-k^2+\frac{N^2k^2}{\sigma^2}-\frac{1}{4H^2}$$

Solutions for acoustic-gravity wave and internal gravity wave exist.

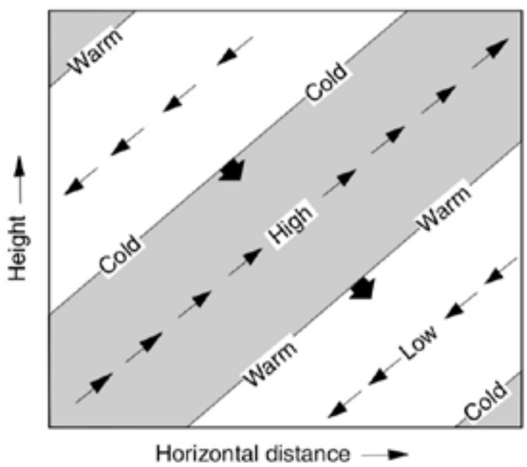
Approximate solution for internal gravity wave is

$$\sigma^2=\frac{N^2k^2}{k^2+m^2+\dfrac{1}{4H^2}} \qquad \rightarrow \sigma < N$$

Structure of gravity wave

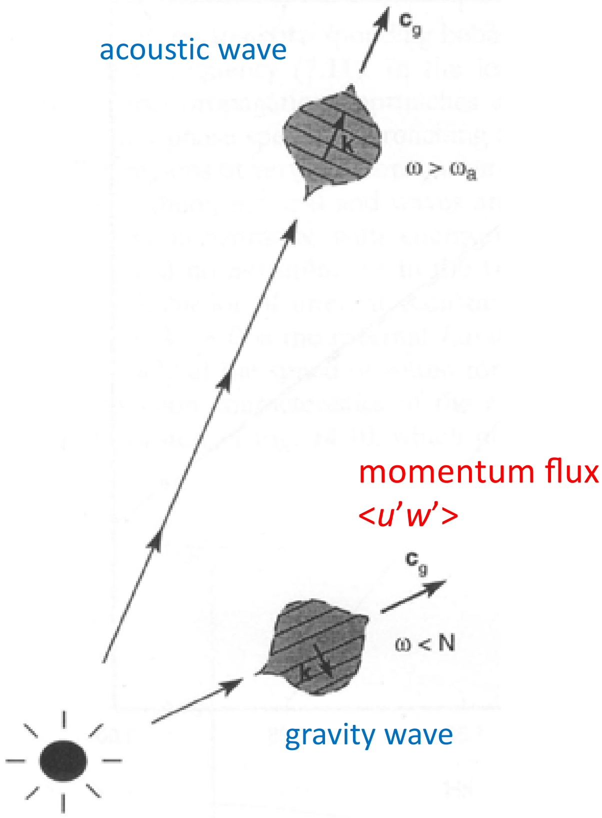
For large-horizontal scale waves (typical in planetary atmospheres),

$$\left|\frac{k}{m}\right|=\left|\frac{\sigma}{N}\right| \qquad \rightarrow \text{Long period waves have near-horizontal phase surfaces}$$

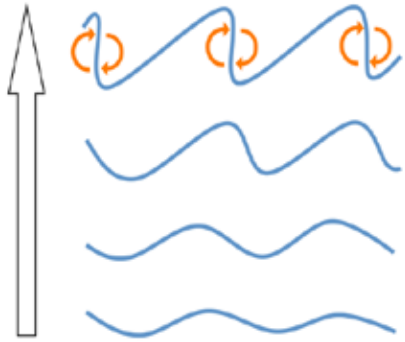


Holton 2004

Fig. 7.9 Idealized cross section showing phases of pressure, temperature, and velocity perturbations for an internal gravity wave. Thin arrows indicate the perturbation velocity field, blunt solid arrows the phase velocity. Shading shows regions of upward motion.

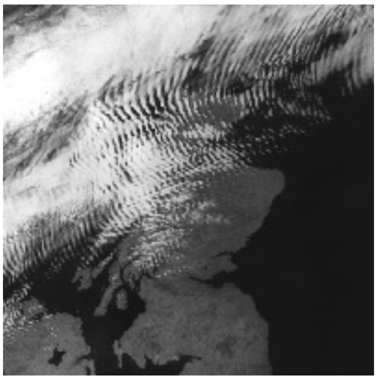


Salby (1996)

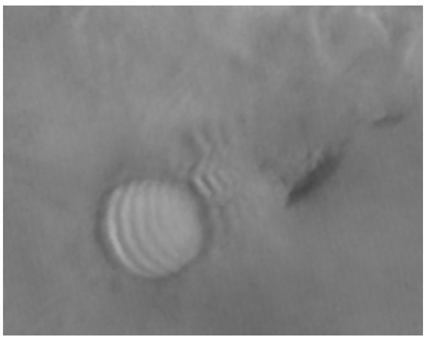


Amplitude growth with height
and wave breaking
→ turbulence generation,
mean-wind acceleration

Observed gravity waves



Mountain waves on Earth



Mountain waves on Mars

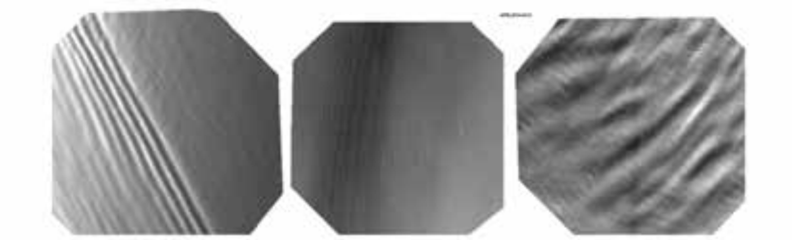
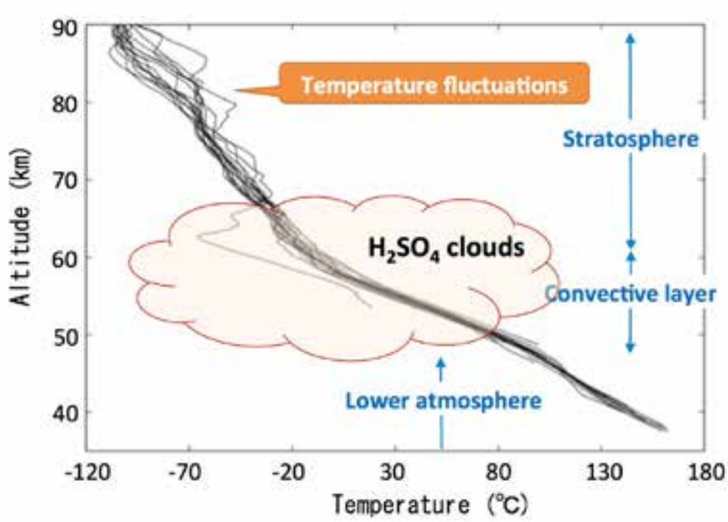


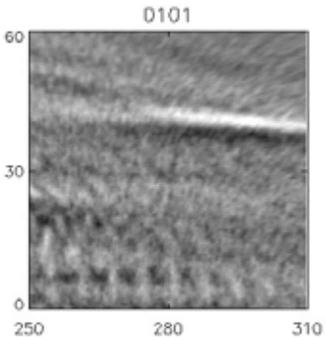
Fig. 11. VMC images of polar waves: left – long waves (NIR1 filter), middle – long waves producing short wave trains (UV), right – irregular waves (UV).

on Venus



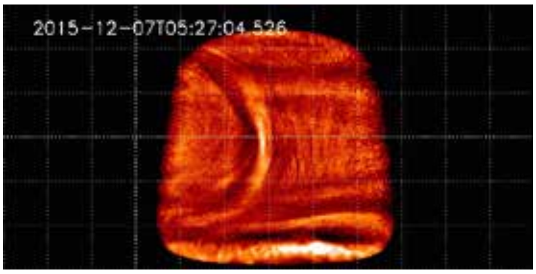
Topographically-generated gravity waves (mountain waves) on Venus

283 nm（mostly SO_2）



Kitahara et al. (2020)

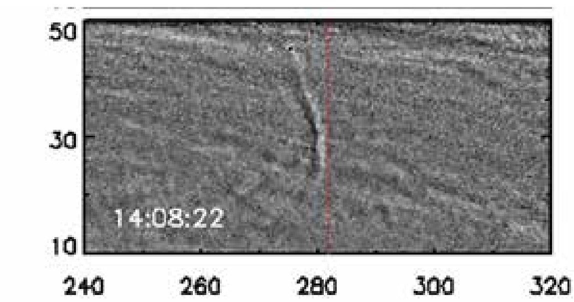
10 μm



Fukuhara et al. 2017
Kouyama et al. 2017

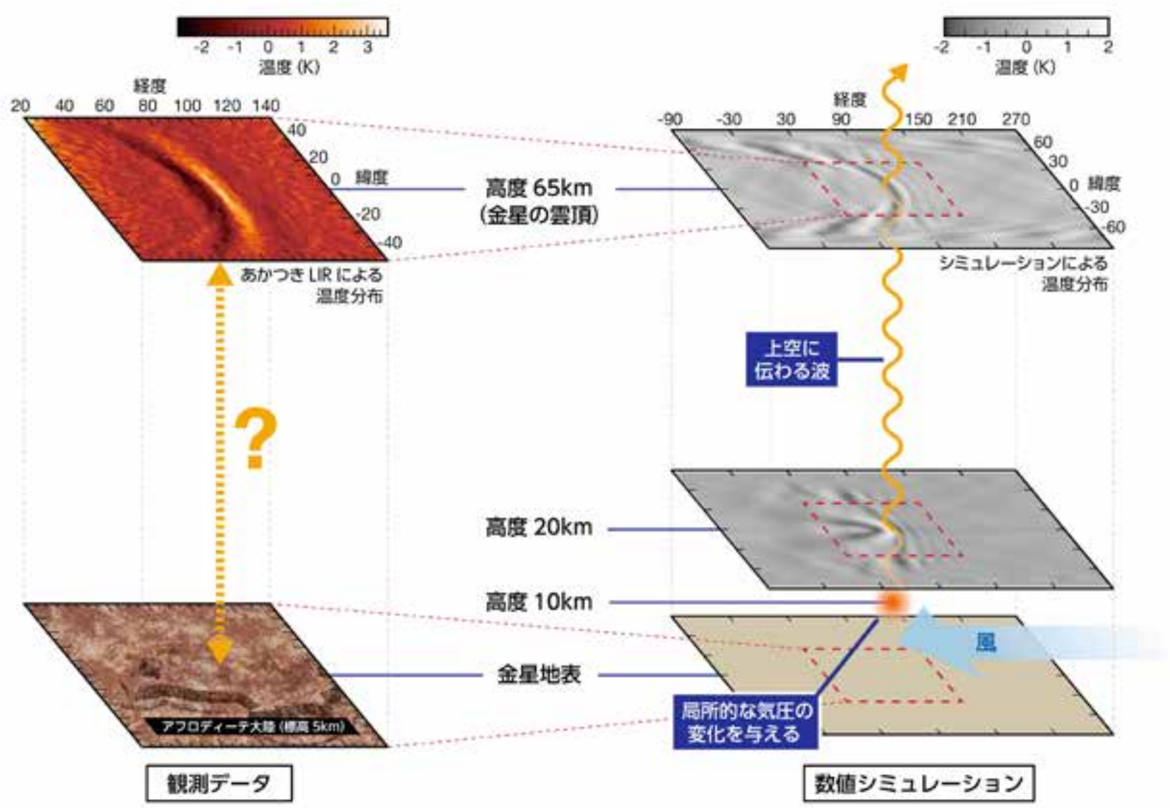
2.02 μm（cloud top height）

Sato et al. (2020)



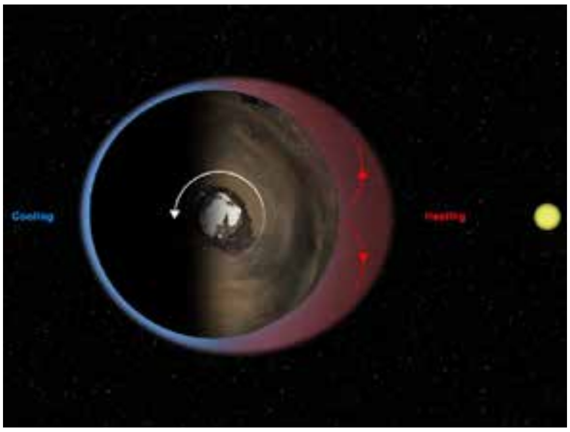
Mountain waves appear mostly in the afternoon region

Mountain wave discovered on Venus



Thermal tides

Planetary-scale gravity wave generated by the movement of the solar heating region in the diurnal cycle



Excitation mechanism:

- Earth : solar heating of stratospheric ozone layer
- Venus : solar heating of cloud layer
- Mars : solar heating of atmospheric dust

Thermal tide in Earth’s atmosphere

Vertical structure

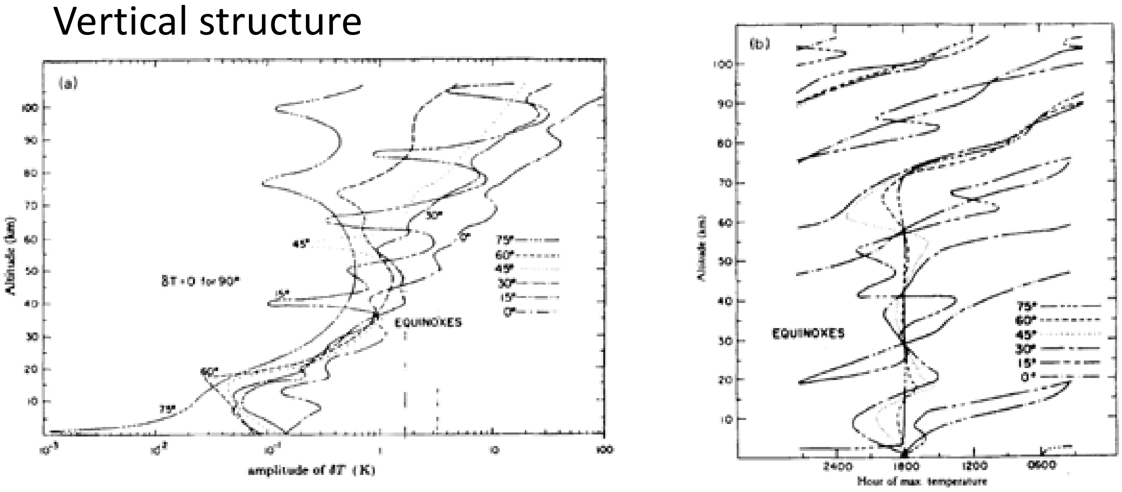


Fig. 4.7. (a) Amplitude and (b) phase of solar diurnal component of T' at various latitudes for equinox. [After Lindzen (1967).]

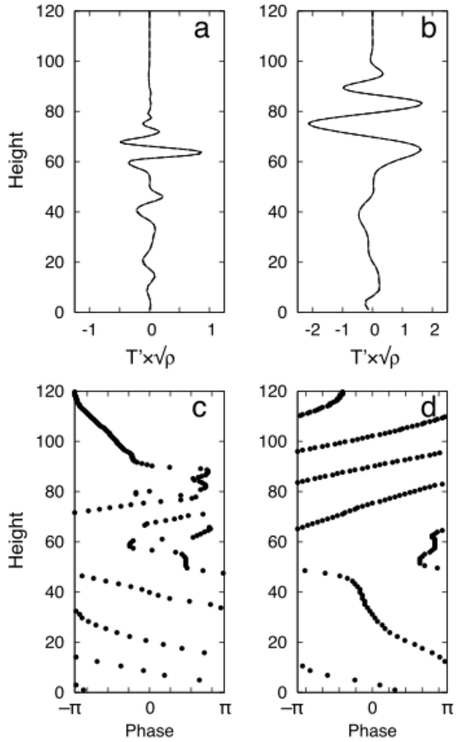
Holton (1992)

Thermal tide in Venus's atmosphere

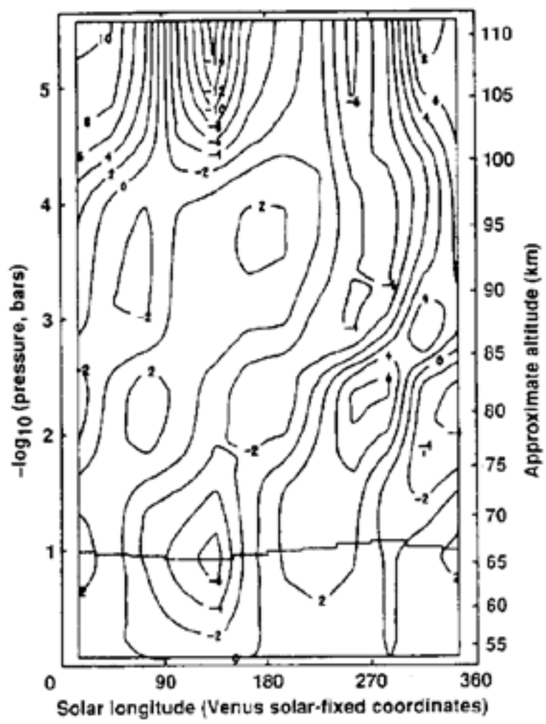
Linear solution (Takagi & Matsuda, 2006)

Dirunal

Semi-diurnal



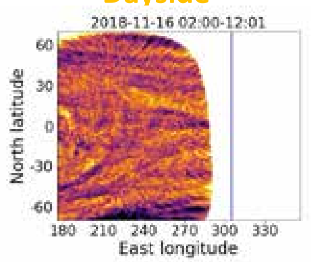
Temperature perturbation
(Schofield & Taylor 1983)



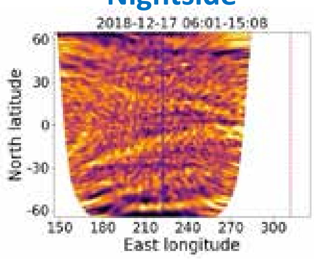
Wind field of thermal tides on Venus observed by Akatsuki infrared camera (Fukuya et al. 2021)

Clouds seen in thermal IR

Dayside



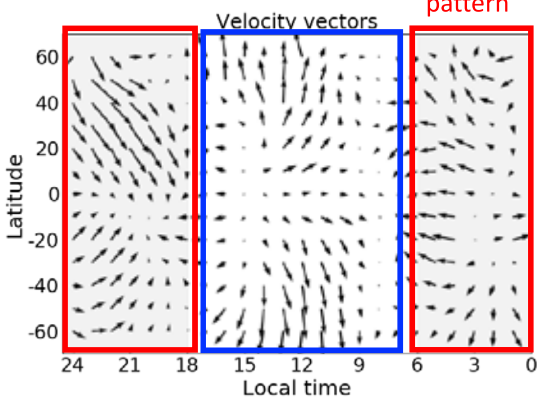
Nightside

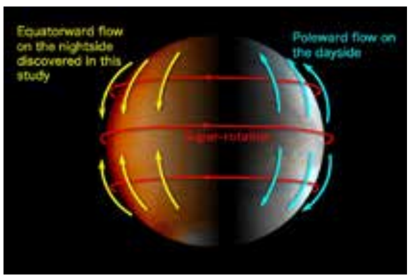


Equatorward

Poleward

Disturbed pattern





- Equatorward circulation was discovered on the nightside
- Dayside poleward flow and the nightside equatorward flow cancel each other